## Linear State-Variable Equations

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## Contents

1 Introduction ..... 2
2 Solving Linear State-Variable Equations ..... 6
2.1 Explicit Solution ..... 6
2.2 Laplace Transform ..... 8
3 Matrix Computations ..... 9
3.1 Determinant ..... 10
Example 1 (Determinant of a Scalar) ..... 10
Example $2 \quad$ (Determinant $M \in \mathbb{R}^{2 \times 2}$ ) ..... 10
Example 3 (Determinant $M \in \mathbb{R}^{3 \times 3}$ ) ..... 11
3.2 Matrix Inverse ..... 12
Example 4 (Inverse $M \in \mathbb{R}^{2 \times 2}$ ) ..... 13
Example 5 (Free response of LTI System $A, B, C, D$ ) ..... 13

## 1 Introduction

The state variable equations are generally in the form

$$
\begin{align*}
& \dot{q}(t)=f(q, u, t)  \tag{1}\\
& y(t)=g(q, u, t) \tag{2}
\end{align*}
$$

where

- $q$ is an $n$-dimensional vector of state variables,
- $u$ is an $p$-dimensional vector of input variables,
- $y$ is a $m$-dimensional vector of output variables,
- $t$ is time, the independent variable of integration, and
- $f$ and $g$ are vector-valued functions of dimension $n$ and $p$ respectively.

This form has several advantages, especially for numerical computation, which is why we learn to derive them from physics-based ODE models.

When the functions $f$ and $g$ are linear functions of the state $q$ and input $u$, and independent of time, the state-variable equations represent a linear time-invariant (LTI) system, for which there are many advanced methods of analysis and design.

For the (LTI) system, we get ODEs

$$
\begin{aligned}
& \dot{q}_{1}=a_{11} q_{1}+a_{12} q_{2}+\cdots a_{1 n} q_{n}+b_{11} u_{1}+b_{12} u_{2}+\cdots b_{1 p} u_{p} \\
& \dot{q}_{2}=a_{21} q_{1}+a_{22} q_{2}+\cdots a_{2 n} q_{n}+b_{21} u_{1}+b_{22} u_{2}+\cdots b_{2 p} u_{p} \\
& \quad \vdots \\
& \dot{q}_{n}=a_{n 1} q_{1}+a_{n 2} q_{2}+\cdots a_{n n} q_{n}+b_{n 1} u_{1}+b_{n 2} u_{2}+\cdots b_{n p} u_{p},
\end{aligned}
$$

and output equations

$$
\begin{aligned}
& y_{1}=c_{11} q_{1}+c_{12} q_{2}+\cdots c_{1 n} q_{n}+d_{11} u_{1}+d_{12} u_{2}+\cdots d_{1 p} u_{p} \\
& y_{2}=c_{21} q_{1}+c_{22} q_{2}+\cdots c_{2 n} q_{n}+d_{21} u_{1}+d_{22} u_{2}+\cdots d_{2 p} u_{p} \\
& \quad \vdots \\
& y_{m}=c_{m 1} q_{1}+c_{m 2} q_{2}+\cdots c_{m n} q_{n}+d_{m 1} u_{1},+d_{m 2} u_{2}+\cdots d_{m p} u_{p}
\end{aligned}
$$

where the coefficients $a_{i j}, b_{i j}, c_{i j}$, and $d_{i j}$ are constants.
The state vector $q$ consists of the $n$ state variables $q_{1}, q_{2}, \ldots, q_{n}$ as follows:

$$
q=\left[\begin{array}{c}
q_{1}  \tag{3}\\
q_{2} \\
\vdots \\
q_{n}
\end{array}\right] .
$$

Table of Contents

Similarly, the state velocity vector $\dot{q}$, initial condition $q(0)$, input $u$ and output $y$ are

$$
\dot{q}=\left[\begin{array}{c}
\dot{q}_{1} \\
\dot{q}_{2} \\
\vdots \\
\dot{q}_{n}
\end{array}\right], \quad q(0)=\left[\begin{array}{c}
q_{1}(0) \\
q_{2}(0) \\
\vdots \\
q_{n}(0)
\end{array}\right], \quad u=\left[\begin{array}{c}
u_{1} \\
u_{2} \\
\vdots \\
u_{m}
\end{array}\right], \quad y=\left[\begin{array}{c}
\dot{y}_{1} \\
\dot{y}_{2} \\
\vdots \\
\dot{y}_{n}
\end{array}\right] .
$$

We can collect the coefficients $a_{i j}, b_{i j}, c_{i j}$, and $d_{i j}$ into matrices as

$$
\left.\begin{array}{rl}
A_{n \times n} & =\left[\begin{array}{cccc}
a_{11} & a_{12} & \cdots & a_{1 n} \\
a_{21} & a_{22} & \cdots & a_{2 n} \\
\vdots & \vdots & \ddots & \vdots \\
a_{n 1} & a_{n 2} & \cdots & a_{n n}
\end{array}\right], \quad B_{n \times p}=\left[\begin{array}{ccc}
b_{11} & b_{12} & \cdots
\end{array} b_{1 p}\right. \\
b_{21} & b_{22} \\
\cdots & b_{2 p} \\
\vdots & \vdots \\
\ddots & \vdots \\
b_{n 1} & b_{n 2}
\end{array} \cdots, b_{n p}\right] ~\left[\begin{array}{cccc}
c_{11} & c_{12} & \cdots & c_{1 n} \\
c_{21} & c_{22} & \cdots & c_{2 n} \\
\vdots & \vdots & \ddots & \vdots \\
c_{m 1} & c_{m 2} & \cdots & c_{m n}
\end{array}\right], \quad D_{m \times p}=\left[\begin{array}{cccc}
d_{11} & d_{12} & \cdots & d_{1 p} \\
d_{21} & d_{22} & \cdots & d_{2 p} \\
\vdots & \vdots & \ddots & \vdots \\
C_{m 1} & d_{m 2} & \cdots & d_{m p}
\end{array}\right] .
$$

and then rewrite the state-variable equations, also known as the state-space equations, as

$$
\begin{gather*}
\dot{q}=A q+B u  \tag{4}\\
y=C x+D u \tag{5}
\end{gather*}
$$

Problem 1 (Textbook Example 3.14). Write the state-variable equations for the problem in Example 3.7 with output given by $f_{s_{1}}$.


Solution (includes solution to Example 3.7):
We draw the free-body diagrams:



Applying Newton's second law, we get

$$
\begin{align*}
m \ddot{q}_{1} & =f_{a}-k_{1}\left(q_{1}-q_{2}\right)-c_{1} \dot{q}_{1}, \text { and }  \tag{6}\\
0 & =k_{1}\left(q_{1}-q_{2}\right)-c_{2} \dot{q}_{2}-k_{2} q_{2} . \tag{7}
\end{align*}
$$

How do we choose the state? We explored two concepts:

1. The state should allow the output to be predicted/computed
2. The state may allow prediction of the positions and/or velocities of all objects in an inertial reference frame.

So, since $f_{s_{1}}=k_{1}\left(q_{1}-q_{2}\right)$, we could again consider a state $q_{R}=q_{1}-q_{2}$, and $v_{R}=\dot{q}_{R}$. We can write the equations

$$
\begin{align*}
\dot{q}_{R} & =v_{R}  \tag{8}\\
\dot{v}_{R} & =\ddot{q}_{1}-\ddot{q}_{2}  \tag{9}\\
& =\frac{1}{m_{1}}\left(f_{a}-k_{1} q_{R}-c_{1} \dot{q}_{1}\right)-\ddot{q}_{2} \tag{10}
\end{align*}
$$

Why won't this work? Ans: because of the $\ddot{q}_{2}$ term.
Idea: differentiate (7) to get an expression for $\ddot{q}_{2}$ in terms of $v_{1}, v_{2}$. The expression would be

$$
\dot{v}_{2}=\frac{1}{c_{2}}\left(k_{1} v_{R}-k_{2} v_{2}\right) .
$$

Then, we would need to add $v_{2}$ as a state. We don't need to add $v_{1}$ as a state, because $v_{1}=v_{2}+v_{R}$.

We finally get

$$
\begin{align*}
\dot{q}_{R} & =v_{R}  \tag{11}\\
\dot{v}_{R} & =\frac{1}{m}\left(f_{a}-k_{1} q_{R}-c_{1}\left(v_{2}+v_{R}\right)\right)-\frac{1}{c_{2}}\left(k_{1} v_{R}-k_{2} v_{2}\right)  \tag{12}\\
\dot{v}_{2} & =\frac{1}{c_{2}}\left(k_{1} v_{R}-k_{2} v_{2}\right)  \tag{13}\\
y & =k_{1} q_{R} \tag{14}
\end{align*}
$$

To get to the matrix form, we expand the terms:

$$
\begin{align*}
\dot{q}_{R} & =v_{R}  \tag{15}\\
\dot{v}_{R} & =-\frac{k_{1}}{m} q_{R}-\frac{c_{1}}{m} v_{R}-\frac{c_{1}}{m} v_{2}+\frac{1}{m} f_{a}  \tag{16}\\
\dot{v}_{2} & =\frac{k_{1}}{c_{2}} v_{R}-\frac{k_{2}}{c_{2}} v_{2}  \tag{17}\\
y & =k_{1} q_{R} \tag{18}
\end{align*}
$$

An intermediate step that might make things easier:

$$
\begin{align*}
\dot{q}_{R} & =0 q_{r}+v_{R}+0 v_{2}+0 f_{a}  \tag{19}\\
\dot{v}_{R} & =-\frac{k_{1}}{m} q_{R}-\frac{c_{1}}{m} v_{R}-\frac{c_{1}}{m} v_{2}+\frac{1}{m} f_{a}  \tag{20}\\
\dot{v}_{2} & =0 q_{R}+\frac{k_{1}}{c_{2}} v_{R}-\frac{k_{2}}{c_{2}} v_{2}+0 f_{a}  \tag{21}\\
y & =k_{1} q_{R}+0 v_{R}+0 v_{2}+0 f_{a} \tag{22}
\end{align*}
$$

With state vector $x=\left[\begin{array}{lll}q_{R} & v_{R} & v_{2}\end{array}\right]^{T}$, we may write

$$
\begin{align*}
& \dot{x}=\left[\begin{array}{ccc}
0 & 1 & 0 \\
-\frac{k_{1}}{m} & -\frac{c_{1}}{m} & -\frac{c_{1}}{m} \\
0 & \frac{k_{1}}{c_{2}} & \frac{k_{2}}{c_{2}}
\end{array}\right] x+\left[\begin{array}{c}
0 \\
\frac{1}{m} \\
0
\end{array}\right] f_{a}  \tag{23}\\
& y=\left[\begin{array}{lll}
k_{1} & 0 & 0
\end{array}\right] x \tag{24}
\end{align*}
$$

Instead, consider the textbook solution, which uses the state $x=\left[\begin{array}{lll}q_{1} & v_{1} & q_{2}\end{array}\right]^{T}$ The statevariable (or state-space) equations are :

$$
\begin{align*}
\dot{q}_{1} & =v_{1}  \tag{25}\\
\dot{v}_{1} & =\frac{1}{m}\left(f_{a}-k_{1}\left(q_{1}-q_{2}\right)-c_{1} v_{1}\right)  \tag{26}\\
\dot{q}_{2} & =\frac{k_{1}}{c_{2}}\left(q_{1}-q_{2}\right)-\frac{k_{2}}{c_{2}} q_{2}  \tag{27}\\
y & =k_{1}\left(q_{1}-q_{2}\right) \tag{28}
\end{align*}
$$

An intermediate step that might make things easier:

$$
\begin{align*}
\dot{q}_{1} & =0 q_{1}+v_{1}+0 q_{1}+0 f_{a}  \tag{29}\\
\dot{v}_{1} & =-\frac{k_{1}}{m} q_{1}-\frac{c_{1}}{m} v_{1}+\frac{k_{1}}{m} q_{2}+\frac{1}{m} f_{a}  \tag{30}\\
\dot{q}_{2} & =\frac{k_{1}}{c_{2}} q_{1}+0 v_{1}-\frac{\left(k_{1}+k_{2}\right)}{c_{2}} q_{2}+0 f_{a}  \tag{31}\\
y & =k_{1} q_{1}+0 v_{1}-k_{1} q_{2}+0 f_{a} \tag{32}
\end{align*}
$$

With state vector, we may write

$$
\begin{align*}
& \dot{x}=\left[\begin{array}{ccc}
0 & 1 & 0 \\
-\frac{k_{1}}{m} & -\frac{c_{1}}{m} & \frac{k_{1}}{m} \\
\frac{k_{1}}{c_{2}} & 0 & -\frac{k_{1}+k_{2}}{c_{2}}
\end{array}\right] x+\left[\begin{array}{c}
0 \\
\frac{1}{m} \\
0
\end{array}\right] f_{a}  \tag{33}\\
& y=\left[\begin{array}{lll}
k_{1} & 0 & -k_{1}
\end{array}\right] x \tag{34}
\end{align*}
$$

## 2 Solving Linear State-Variable Equations

Suppose that we are given the state-variable equations

$$
\begin{align*}
& \dot{x}(t)=A x(t)+B u(t)  \tag{1}\\
& y(t)=C x(t)+D u(t) \tag{2}
\end{align*}
$$

where $x(t)$ is the state, $y(t)$ is the output, $u(t)$ is the input, and $t$ represents time.
Suppose we know the initial condition (IC) $x\left(t_{0}\right)$, and input $u(t)$ for $t \in\left[t_{0}, t_{f}\right]$. We want to understand how the output $y(t)$ will behave over the time interval $\left[t_{0}, t_{f}\right]$. To do so, we may either

- Explicitly solve for $x(t)$, because $y(t)=C x(t)+D u(t)$
- Use $A, B, C$, and $D$ to predict the behavior of solutions $x(t)$ given ICs and input.

We saw that transfer functions allow us to do something similar for Input-Output Differential Equations.

### 2.1 Explicit Solution

In earlier calculus classes, you may have seen methods to solve linear ODEs by computing homogenous and particular solutions. In the notes on Laplace transforms, we solve firstorder input-output differential equations in $y(t)$ using this method. This section shows the relationship between that method and the linear state-variable equations given by matrices $A, B, C$, and $D$.

Matrix Exponential. Given a matrix $A$, we define the matrix $e^{A t}$ as the infinite sequence

$$
\begin{equation*}
e^{A t}=I+\frac{1}{1} A t+\frac{1}{2!} A^{2} t^{2}+\frac{1}{3!} A^{3} t^{3}+\frac{1}{4!} A^{4} t^{4}+\cdots \tag{3}
\end{equation*}
$$

This definition implies that $e^{A\left(t_{1}+t_{2}\right)}=e^{A t_{1}} e^{A t_{2}}$.
Let's calculate the derivative of $e^{A t}$ :

$$
\begin{aligned}
\frac{d}{d t} e^{A t} & =\frac{d}{d t}\left(I+\frac{1}{1} A t+\frac{1}{2!} A^{2} t^{2}+\frac{1}{3!} A^{3} t^{3}+\frac{1}{4!} A^{4} t^{4}+\cdots\right) \\
& =\frac{d}{d t}(I)+\frac{d}{d t}\left(\frac{1}{1} A t\right)+\frac{d}{d t}\left(\frac{1}{2!} A^{2} t^{2}\right)+\frac{d}{d t}\left(\frac{1}{3!} A^{3} t^{3}\right)+\frac{d}{d t}\left(\frac{1}{4!} A^{4} t^{4}\right)+\cdots \\
& =0+A+\frac{1}{2!} A^{2}(2 t)+\frac{1}{3!} A^{3}\left(3 t^{2}\right)+\frac{1}{4!} A^{4}\left(4 t^{3}\right)+\cdots \\
& =0+A+\frac{1}{1} A^{2} t+\frac{1}{2!} A^{3} t^{2}+\frac{1}{3!} A^{4} t^{3}+\cdots \\
& =A\left(I+\frac{1}{1} A t+\frac{1}{2!} A^{2} t^{2}+\frac{1}{3!} A^{3} t^{3}+\cdots\right) \\
& =A e^{A t}
\end{aligned}
$$

Suppose we define $z(t)=e^{A t} v$, where $v=e^{-A t_{0}} z\left(t_{0}\right)$, Then,

$$
\begin{aligned}
\dot{z}(t) & =\frac{d}{d t}\left(e^{A t} v\right)=\frac{d}{d t}\left(e^{A t}\right) v \\
& =\left(A e^{A t}\right) v=A\left(e^{A t} v\right) \\
& =A z(t)
\end{aligned}
$$

In other words, $x(t)=e^{A\left(t-t_{0}\right)} x\left(t_{0}\right)$ is the solution to the differential equation

$$
\dot{x}(t)=A x(t)
$$

with initial condition $x\left(t_{0}\right)$.
Suppose we have the differential equation

$$
\dot{x}(t)=A x(t)+B u(t)
$$

Multiply on the left by $e^{-A t}$ and rearrange to get

$$
e^{-A t} \dot{x}(t)-A e^{-A t} x(t)=e^{-A t} B u(t)
$$

But,

$$
\frac{d}{d t}\left(e^{-A t} x(t)\right)=e^{-A t} \dot{x}(t)-A e^{-A t} x(t)
$$

so that we may write

$$
\frac{d}{d t}\left(e^{-A t} x(t)\right)=e^{-A t} B u(t)
$$

Table of Contents

Integrate this equation on both sides:

$$
\begin{aligned}
& \int_{t_{0}}^{t} \frac{d}{d t}\left(e^{-A \tau} x(\tau)\right) d \tau=\int_{t_{0}}^{t} e^{-A \tau} B u(\tau) d \tau \\
\Longrightarrow & e^{-A t} x(t)-e^{-A t_{0}} x\left(t_{0}\right)=\int_{t_{0}}^{t} e^{-A \tau} B u(\tau) d \tau
\end{aligned}
$$

Now, multiply both sides by $e^{A t}$ :

$$
\begin{gathered}
e^{A t} e^{-A t} x(t)-e^{A t} e^{-A t_{0}} x\left(t_{0}\right)=e^{A t} \int_{t_{0}}^{t} e^{-A \tau} B u(\tau) d \tau \\
\Longrightarrow x(t)-e^{A\left(t-t_{0}\right)} x\left(t_{0}\right)=\int_{t_{0}}^{t} e^{A(t-\tau)} B u(\tau) d \tau \\
\Longrightarrow x(t)=e^{A\left(t-t_{0}\right)} x\left(t_{0}\right)+\int_{t_{0}}^{t} e^{A(t-\tau)} B u(\tau) d \tau \\
\Longrightarrow y(t)=C e^{A\left(t-t_{0}\right)} x\left(t_{0}\right)+\int_{t_{0}}^{t} C e^{A(t-\tau)} B u(\tau) d \tau+D u(t) \\
=\text { free response }+ \text { forced response }
\end{gathered}
$$

The forced response is $\int_{t_{0}}^{t} C e^{A(t-\tau)} B u(\tau) d \tau+D u(t)$, which may be difficult to calculate. The first term in the forced response is exactly a convolution operation between the function $C e^{A t}$ and $B u(t)$.

### 2.2 Laplace Transform

Again, if the goal is to explicitly calculate $y(t)$, we may prefer to work in the $s$-domain, which implies we work with Laplace transforms.

$$
\begin{gather*}
\dot{x}(t)=A x(t)+B u(t)  \tag{4}\\
y(t)=C x(t)+D u(t)  \tag{5}\\
\Longrightarrow \begin{array}{l}
s \hat{x}(s)-x\left(t_{0}\right)
\end{array}=A \hat{x}+B \hat{u}(s)  \tag{6}\\
\Longrightarrow(s I-A) \hat{x}(s)=x\left(t_{0}\right)+B \hat{u}(s)  \tag{7}\\
\hat{x}(s)=(s I-A)^{-1} x\left(t_{0}\right)+(s I-A)^{-1} B \hat{u}(s)  \tag{8}\\
 \tag{9}\\
\Longrightarrow \hat{y}(s)=C \hat{x}(s)+D \hat{u}(s)  \tag{10}\\
\\
\hline \hat{y}(s)=C(s I-A)^{-1} x\left(t_{0}\right)+C(s I-A)^{-1} B \hat{u}(s)+D \hat{u}(s)
\end{gather*}
$$

To find the Laplace transform, set $x\left(t_{0}\right)=0$ to obtain

$$
\begin{align*}
\hat{y}(s) & =C(s I-A)^{-1} B \hat{u}(s)+D \hat{u}(s)  \tag{11}\\
\Longrightarrow \hat{y}(s) & =\left(C(s I-A)^{-1} B+D\right) \hat{u}(s)=G(s) \hat{u}(s)  \tag{12}\\
\Longrightarrow G(s) & =C(s I-A)^{-1} B+D \tag{13}
\end{align*}
$$

The remainder of this subsection shows that the explicit time-domain solution would be the same as if we used the $s$-domain and the inverse Laplace transform to calculate $y(t)$.

Fact. Let $A$ be an $n \times n$ matrix. Then,

$$
\mathcal{L}\left\{e^{A t}\right\}=(s I-A)^{-1} e^{A t_{0}} \Longrightarrow \mathcal{L}^{-1}\left\{(s I-A)^{-1}\right\}=e^{A\left(t-t_{0}\right)}
$$

To derive this fact, apply the Differentiation rule:

$$
\begin{array}{rlr} 
& \mathcal{L}\left\{e^{A t}\right\}=\mathcal{L}\left\{e^{A t}\right\} & \text { (Always true) } \\
\Longrightarrow & s \mathcal{L}\left\{e^{A t}\right\}-e^{A t_{0}}=\mathcal{L}\left\{\frac{d}{d t}\left(e^{A t}\right)\right\} & \text { (Differentiation rule) } \\
\Longrightarrow & s \mathcal{L}\left\{e^{A t}\right\}-e^{A t_{0}}=\mathcal{L}\left\{\left(A e^{A t}\right)\right\} & \text { (derived earlier) } \\
\Longrightarrow & s \mathcal{L}\left\{e^{A t}\right\}-e^{A t_{0}}=A \mathcal{L}\left\{e^{A t}\right\} & \text { (Linearity of LT) } \\
\Longrightarrow & (s I-A) \mathcal{L}\left\{e^{A t}\right\}=e^{A t_{0}} & \text { (Rearrange terms) } \\
\Longrightarrow & \mathcal{L}\left\{e^{A t}\right\}=(s I-A)^{-1} e^{A t_{0}} & \text { (Matrix inversion) }
\end{array}
$$

So, we can take the inverse Laplace transform of (10) to obtain

$$
\begin{aligned}
\mathcal{L}^{-1}\{\hat{y}(s)\} & =\mathcal{L}^{-1}\left\{C(s I-A)^{-1} x\left(t_{0}\right)+C(s I-A)^{-1} B \hat{u}(s)+D \hat{u}(s)\right\} \\
y(t) & =\mathcal{L}^{-1}\left\{C(s I-A)^{-1} x\left(t_{0}\right)\right\}+\mathcal{L}^{-1}\left\{C(s I-A)^{-1} B \hat{u}(s)\right\}+\mathcal{L}^{-1}\{D \hat{u}(s)\} \\
y(t) & =C e^{A\left(t-t_{0}\right)} x\left(t_{0}\right)+\left(\mathcal{L}^{-1}\left\{C(s I-A)^{-1}\right\} * \mathcal{L}^{-1}\{B \hat{u}(s)\}\right)(t)+D \mathcal{L}^{-1}\{\hat{u}(s)\}
\end{aligned}
$$

(Convolution of functions of $t$ is product of Laplace transforms of those functions)

$$
\begin{aligned}
& y(t)=C e^{A\left(t-t_{0}\right)} x\left(t_{0}\right)+\left(C e^{A t} * B u(t)\right)(t)+D u(t) \\
& y(t)=C e^{A t} x(0)+\int_{t_{0}}^{t} C e^{A(t-\tau)} B u(\tau) d \tau+D u(t)
\end{aligned}
$$

## 3 Matrix Computations

To calculate the transfer function, we must compute $(s I-A)^{-1}$, where $M^{-1}$ is the inverse of a square matrix $M$.

To calculate the inverse of $M$, we need to calculate the determinant of $M$.

### 3.1 Determinant

The determinant of $M$, and $n \times n$ matrix, is always a scalar number.
Let $M_{i, j}$ be the $(i, j)^{\text {th }}$ element of $M$.

Let $S_{n}$ be the set of $n$ ! possible permutations of an ordered set of $n$ numbers.

The determinant of a $n \times n$ square matrix $M$ is given by

$$
\operatorname{det} M=\sum_{\sigma \in S_{n}}(-1)^{N_{\sigma}} \Pi_{i}^{n} M_{i, \sigma(i)}
$$

where $N_{\sigma}$ is the number of pairwise exchanges of elements of $\sigma$ required to reach the order $(1,2, \ldots, n)$.

Example 1 (Determinant of a Scalar). Let

$$
M=[a]
$$

Calculate $\operatorname{det} M$.
Solution:
$S_{1}=\{(1)\}=\left\{\sigma_{1}\right\} . N_{\sigma_{1}}=0$, because we don't need to switch any elements to get to the permutation (1). det $M$ has only one term:

$$
\begin{align*}
\operatorname{det} M & =(-1)^{N_{\sigma_{1}}} \Pi_{i}^{1} M_{i, \sigma_{1}(i)}  \tag{1}\\
& =(-1)^{0} M_{1, \sigma_{1}(1)}  \tag{2}\\
& =1 \cdot M_{1,1}  \tag{3}\\
& =a \tag{4}
\end{align*}
$$

Example 2 (Determinant $M \in \mathbb{R}^{2 \times 2}$ ). Let

$$
M=\left[\begin{array}{ll}
a & b \\
c & d
\end{array}\right]
$$

Calculate $\operatorname{det} M$.

Solution: Since $M$ is a $2 \times 2$ matrix, we will need to work with $S_{2}$, which has $2!=2$ elements: $S_{2}=\{(1,2), \quad(2,1)\}=\left\{\sigma_{1}, \sigma_{2}\right\}$.
$\operatorname{det} M$ is the sum of two terms:

$$
\begin{equation*}
\operatorname{det} M=(-1)^{N_{\sigma_{1}}} \Pi_{i}^{2} M_{i, \sigma_{1}(i)}+(-1)^{N_{\sigma_{2}}} \Pi_{i}^{2} M_{i, \sigma_{2}(i)} \tag{5}
\end{equation*}
$$

Consider the first term corresponding to $\sigma_{1}=(1,2) . N_{\sigma_{1}}=0$, because we don't need to permute any entries to reach the permutation (1,2). Since $\sigma_{1}=(1,2)$, we have

$$
\sigma_{1}(1)=1, \sigma_{1}(2)=2
$$

We need to evaluate $(-1)^{N_{\sigma_{1}}} \Pi_{i}^{2} M_{i, \sigma_{1}(i)}$ :

$$
\begin{align*}
(-1)^{N_{\sigma_{1}}} \Pi_{i}^{2} M_{i, \sigma_{1}(i)} & =(-1)^{N_{\sigma_{1}}} M_{1, \sigma_{1}(1)} M_{2, \sigma_{1}(2)}  \tag{6}\\
& =(-1)^{0} M_{1,1} M_{2,2}  \tag{7}\\
& =a b \tag{8}
\end{align*}
$$

For $\sigma_{2}=\{(2,1)\}: N_{\sigma_{2}}=1$, because we must switch the 1 and 2 to obtain the permutation $(1,2)$. We have

$$
\begin{align*}
\sigma_{2}(1) & =2, \sigma_{2}(2)=1 \\
(-1)^{N_{\sigma_{2}}} \Pi_{i}^{2} M_{i, \sigma_{2}(i)} & =(-1)^{N_{\sigma_{2}}} M_{1, \sigma_{2}(1)} M_{2, \sigma_{2}(2)}  \tag{9}\\
& =(-1)^{1} M_{1,2} M_{2,1}  \tag{10}\\
& =-b c \tag{11}
\end{align*}
$$

Therefore, $\operatorname{det} M=a d-b c$

Example 3 (Determinant $M \in \mathbb{R}^{3 \times 3}$ ). Let

$$
M=\left[\begin{array}{lll}
a & b & c \\
d & e & f \\
g & h & i
\end{array}\right]
$$

Calculate $\operatorname{det} M$.
Solution: $M \in \mathbb{R}^{3 \times 3}$. Therefore, the permutations we consider belong to $S_{3}$. There are $3!=6$ such permutations. They are:
$S_{3}=\{(1,2,3), \quad(1,3,2), \quad(2,1,3), \quad(2,3,1), \quad(3,1,2), \quad(3,2,1)\}=\left\{\sigma_{1}, \sigma_{2}, \sigma_{3}, \sigma_{4}, \sigma_{5}, \sigma_{6}\right\}$
$\operatorname{det} M$ is the sum of six terms:

$$
\begin{align*}
\operatorname{det} M= & (-1)^{N_{\sigma_{1}}} \Pi_{i}^{3} M_{i, \sigma_{1}(i)}+(-1)^{N_{\sigma_{2}}} \Pi_{i}^{3} M_{i, \sigma_{2}(i)}+(-1)^{N_{\sigma_{3}}} \Pi_{i}^{3} M_{i, \sigma_{3}(i)}  \tag{12}\\
& +(-1)^{N_{\sigma_{4}}} \Pi_{i}^{3} M_{i, \sigma_{4}(i)}+(-1)^{N_{\sigma_{5}}} \Pi_{i}^{3} M_{i, \sigma_{5}(i)}+(-1)^{N_{\sigma_{6}}} \Pi_{i}^{3} M_{i, \sigma_{6}(i)}
\end{align*}
$$

Consider the first term corresponding to $\sigma_{1}=(1,2,3) . N_{\sigma_{1}}=0$, because we don't need to permute any entries to reach the permutation $(1,2,3)$. Since $\sigma_{1}=(1,2,3)$, we have

$$
\sigma_{1}(1)=1, \sigma_{1}(2)=2, \sigma_{1}(3)=3
$$

We need to evaluate $(-1)^{N_{\sigma_{1}}} \Pi_{i}^{3} M_{i, \sigma_{1}(i)}$ :

$$
\begin{align*}
(-1)^{N_{\sigma_{1}}} \Pi_{i}^{3} M_{i, \sigma_{1}(i)} & =(-1)^{0} M_{1, \sigma_{1}(1)} M_{2, \sigma_{1}(2)} M_{3, \sigma_{1}(3)}  \tag{13}\\
& =M_{1,1} M_{2,2} M_{3,3}  \tag{14}\\
& =a e i \tag{15}
\end{align*}
$$

So, the first term in $\operatorname{det} M$ is $+a e i$.
We repeat this process for the second term in det $M$ corresponding to $\sigma_{2}=(1,3,2) . N_{\sigma_{2}}=1$, because we need to exchange the last two elements of $\sigma_{2}$ to get the permutation $(1,2,3)$. Since $\sigma_{2}=(1,3,2)$, we have

$$
\begin{align*}
\sigma_{2}(1)=1 & , \sigma_{2}(2)=3, \sigma_{2}(3)=2 \\
(-1)^{N_{\sigma_{2}}} \Pi_{i}^{3} M_{i, \sigma_{2}(i)} & =(-1)^{1} M_{1, \sigma_{2}(1)} M_{2, \sigma_{2}(2)} M_{3, \sigma_{2}(3)}  \tag{16}\\
& =-M_{1,1} M_{2,3} M_{3,2}  \tag{17}\\
& =-a f h \tag{18}
\end{align*}
$$

Continuing this process, we get

$$
\operatorname{det} M=a e i-a f h+d h c-d i b+g b f-g c e .
$$

### 3.2 Matrix Inverse

sad Let $M$ be an $n \times n$ matrix.
The inverse of $M$, denoted $M^{-1}$, is a matrix whose $(i, j)^{\text {th }}$ element $M_{i, j}^{-1}$ is given by

$$
M_{i, j}^{-1}=(-1)^{(i+j)} \frac{\operatorname{det} M_{[i, j]}}{\operatorname{det} M}
$$

where $M_{[i, j]}$ is an $(n-1) \times(n-1)$ matrix obtained by deleting the $i^{\text {th }}$ column and $j^{\text {th }}$ row of $M$.

Example 4 (Inverse $M \in \mathbb{R}^{2 \times 2}$ ). Let

$$
M=\left[\begin{array}{ll}
a & b \\
c & d
\end{array}\right]
$$

Calculate $M^{-1}$.

Solution:

$$
\operatorname{det} M=a d-b c
$$

By deleting the $i^{\text {th }}$ column and $j^{\text {th }}$ row of $M$, we get

$$
\begin{align*}
& M_{[1,1]}=d  \tag{19}\\
& M_{[1,2]}=b  \tag{20}\\
& M_{[2,1]}=c  \tag{21}\\
& M_{[2,2]}=a \tag{22}
\end{align*}
$$

The $(i, j)^{\text {th }}$ entry of $M^{-1}$ is then

$$
\begin{align*}
& M_{1,1}^{-1}=(-1)^{(1+1)} \frac{\operatorname{det} M_{[1,1]}}{\operatorname{det} M}=\frac{d}{a d-b c}  \tag{23}\\
& M_{1,2}^{-1}=(-1)^{(1+2)} \frac{\operatorname{det} M_{[1,2]}}{\operatorname{det} M}=\frac{-b}{a d-b c}  \tag{24}\\
& M_{2,1}^{-1}=(-1)^{(2+1)} \frac{\operatorname{det} M_{[2,1]}}{\operatorname{det} M}=\frac{-c}{a d-b c}  \tag{25}\\
& M_{2,2}^{-1}=(-1)^{(2+2)} \frac{\operatorname{det} M_{[2,2]}}{\operatorname{det} M}=\frac{a}{a d-b c} \tag{26}
\end{align*}
$$

Therefore,

$$
M^{-1}=\frac{1}{a d-b c}\left[\begin{array}{cc}
d & -b \\
-c & a
\end{array}\right] .
$$

Example 5 (Free response of LTI System $A, B, C, D$ ).

$$
A=\left[\begin{array}{cc}
-1 & 1  \tag{27}\\
0 & -2
\end{array}\right], C=\left[\begin{array}{ll}
1 & 0
\end{array}\right], B=0, D=0
$$

Let

$$
x(0)=\left[\begin{array}{l}
0 \\
1
\end{array}\right]
$$

Find the free response using Laplace transforms.
Note: $\mathcal{L}\left\{e^{A t}\right\}=(s I-A)^{-1} e^{A t_{0}}$

## Solution:

The free response $\hat{y}_{\text {free }}(s)$ is

$$
\hat{y}_{\text {free }}(s)=C(s I-A)^{-1} x\left(t_{0}\right)
$$

Let's first construct $M=(s I-A)$ :

$$
\begin{aligned}
M=(s I-A) & =s\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right]-\left[\begin{array}{cc}
-1 & 1 \\
0 & -2
\end{array}\right] \\
& =\left[\begin{array}{ll}
s & 0 \\
0 & s
\end{array}\right]-\left[\begin{array}{cc}
-1 & 1 \\
0 & -2
\end{array}\right] \\
& =\left[\begin{array}{cc}
(s+1) & -1 \\
0 & (s+2)
\end{array}\right]
\end{aligned}
$$

Calculate the determinant:

$$
\begin{align*}
\operatorname{det} M & =(s+1)(s+2)-(-1) \cdot(0)  \tag{28}\\
& =(s+1)(s+2) \tag{29}
\end{align*}
$$

We've derived the expression for the inverse of a $2 \times 2$ matrix, so that

$$
\begin{align*}
M^{-1}=(s I-A)^{-1} & =\frac{1}{(s+1)(s+2)}\left[\begin{array}{cc}
(s+2) & 1 \\
0 & (s+1)
\end{array}\right]  \tag{30}\\
& =\left[\begin{array}{cc}
\frac{1}{s+1} & \frac{1}{(s+1)(s+2)} \\
0 & \frac{1}{s+2}
\end{array}\right]  \tag{31}\\
\hat{y}_{\text {free }}(s) & =C(s I-A)^{-1} x\left(t_{0}\right) \\
& =\left[\begin{array}{ll}
1 & 0
\end{array}\right]\left[\begin{array}{cc}
\frac{1}{s+1} & \frac{1}{(s+1)(s+2)} \\
0 & \frac{1}{s+2}
\end{array}\right]\left[\begin{array}{l}
0 \\
1
\end{array}\right] \\
& =\frac{1}{(s+1)(s+2)} \\
& =\frac{1}{s+1}-\frac{1}{s+2} \\
\Longrightarrow y_{\text {free }}(t) & =\mathcal{L}^{-1}\left\{\frac{1}{s+1}\right\}-\mathcal{L}^{-1}\left\{\frac{1}{s+2}\right\} \\
& =e^{-t}-e^{-2 t}
\end{align*}
$$

